

# Fourier Optics Formulas

## Two-Dimensional Fourier Analysis

### Fourier Transform

$$\begin{aligned}\mathcal{F}\{g(x, y)\} &= \iint_{-\infty}^{+\infty} g(x, y) e^{-2\pi(f_X x + f_Y y)} dx dy \\ \mathcal{F}^{-1}\{G(f_X, f_Y)\} &= \iint_{-\infty}^{+\infty} G(f_X, f_Y) e^{2\pi(f_X x + f_Y y)} df_X df_Y \\ \mathcal{F}\{g(ax, by)\} &= \frac{1}{|ab|} G\left(\frac{f_X}{a}, \frac{f_Y}{b}\right)\end{aligned}$$

$$\begin{aligned}\mathcal{F}\{g(x-a, x-b)\} &= G(f_X, f_Y) e^{-j2\pi(f_X a + f_Y b)} \\ \mathcal{F}\{g \star g\} &= |G(f_X, f_Y)|^2\end{aligned}$$

### Fourier-Bessel Transform

$$\begin{aligned}\mathcal{B}\{g_R(r)\} &= 2\pi \int_0^{\infty} r g_R(r) J_0(2\pi r \rho) dr \\ \mathcal{B}^{-1}\{G_0(\rho)\} &= 2\pi \int_0^{\infty} \rho G_0(\rho) J_0(2\pi r \rho) d\rho \\ \mathcal{B}\{g_R(ar)\} &= \frac{1}{a^2} G\left(\frac{\rho}{a}\right)\end{aligned}$$

## Scalar Diffraction Theory

$$u(\mathbf{r}, t) = \operatorname{Re} \{U(\mathbf{r}) e^{-j2\pi\nu t}\}$$

$$(\nabla^2 + k^2)U = 0$$

Huygens-Fresnel Principle (First Rayleigh-Sommerfeld Solution)

$$\begin{aligned}U(\mathbf{r}_0) &= \frac{1}{j\lambda} \iint_{\Sigma} U(\mathbf{r}_1) \frac{\exp(jkr_{01})}{r_{01}} \cos\theta ds \\ U(x, y) &= \frac{z}{j\lambda} \iint_{-\infty}^{+\infty} U(\xi, \eta) \frac{\exp(jkr_{01})}{r_{01}^2} d\xi d\eta\end{aligned}$$

## Angular Spectrum

$$\begin{aligned}A\left(\frac{\alpha}{\lambda}, \frac{\beta}{\lambda}; 0\right) &= \mathcal{F}\{U(x, y, 0)\} \Big|_{f_X = \frac{\alpha}{\lambda}, f_Y = \frac{\beta}{\lambda}} \\ H_z(f_X, f_Y) &= \begin{cases} \exp\left[j2\pi\frac{z}{\lambda}\sqrt{1 - (\lambda f_X)^2 - (\lambda f_Y)^2}\right] & \rho < \frac{1}{\lambda} \\ 0 & \text{otherwise} \end{cases}\end{aligned}$$

## Fresnel (Paraxial) Approximation

$$\begin{aligned}U(x, y) &= \frac{e^{j kz}}{j\lambda z} \iint_{-\infty}^{+\infty} U(\xi, \eta) \exp\left\{j\frac{k}{2z}[(x - \xi)^2 + (y - \eta)^2]\right\} d\xi d\eta \\ U(x, y) &= \frac{e^{j kz}}{j\lambda z} e^{j\frac{k}{2z}(x^2 + y^2)} \iint_{-\infty}^{+\infty} \left[U(\xi, \eta) e^{j\frac{k}{2z}(\xi^2 + \eta^2)}\right] e^{-j\frac{2\pi}{\lambda z}(x\xi + y\eta)} d\xi d\eta\end{aligned}$$

$$H_z(f_X, f_Y) = e^{j kz} \exp[-j\pi\lambda z(f_X^2 + f_Y^2)], \quad |\lambda f| \ll 1$$

## Fraunhofer Approximation

$$\begin{aligned}U(x, y) &= \frac{e^{j kz}}{j\lambda z} e^{j\frac{k}{2z}(x^2 + y^2)} \iint_{-\infty}^{+\infty} U(\xi, \eta) e^{-j\frac{2\pi}{\lambda z}(x\xi + y\eta)} d\xi d\eta \\ U(x, y) &= \frac{e^{j kz}}{j\lambda z} e^{j\frac{k}{2z}(x^2 + y^2)} \mathcal{F}\{U(\xi, \eta)\} \Big|_{f_X = \frac{x}{\lambda z}, f_Y = \frac{y}{\lambda z}}\end{aligned}$$

## Coherent Optical Systems

$$t_l(x, y) = \exp\left[-j\frac{k}{2f}(x^2 + y^2)\right]$$

## Frequency Analysis

### Coherent (Linear in Amplitude)

$$H(f_X, f_Y) = P(\lambda z_i f_X, \lambda z_i f_Y)$$

$$I_i(u, v) = |h * U_g|^2 \subset [H \cdot \mathcal{F}\{U_g\}] \star [H \cdot \mathcal{F}\{U_g\}]$$

Incoherent (Linear in Intensity)

$$\mathcal{H}(f_X, f_Y) = \frac{H \star H}{H \star H \Big|_{f_X=0, f_Y=0}}$$

$$I_i(u, v) = |h|^2 * |U_g|^2 \subset [H \star H][U_g \star U_g]$$

$$\delta = 0.61 \frac{\lambda}{\text{NA}}$$

### Jones' Calculus

$$\mathbf{L}_{\text{rotate}}(\theta) = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

$$\mathbf{L}_{\text{retard}}(\Delta) = \begin{bmatrix} 1 & 0 \\ 0 & e^{-j\Delta} \end{bmatrix}$$

$$\mathbf{L}_{\text{polarizer}}(\alpha) = \begin{bmatrix} \cos^2 \alpha & \sin \alpha \cos \alpha \\ \sin \alpha \cos \alpha & \sin^2 \alpha \end{bmatrix}$$

$$\mathbf{R} = \begin{bmatrix} -1 & 0 \\ 1 & 0 \end{bmatrix}$$

$$\mathbf{L}_{\text{reflect}} = \mathbf{R} \mathbf{L}^t \mathbf{L}$$